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**Astronomy  
&  
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# CARMENES input catalogue of M dwarfs

## II. High-resolution imaging with FastCam<sup>★</sup>

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### ABSTRACT

**Aims.** We search for low-mass companions of M dwarfs and characterize their multiplicity fraction with the purpose of helping in the selection of the most appropriate targets for the CARMENES exoplanet survey.

**Methods.** We obtained high-resolution images in the *I* band with the lucky imaging instrument FastCam at the 1.5 m Telescopio Carlos Sánchez for 490 mid- to late-M dwarfs. For all the detected binaries, we measured angular separations, position angles, and magnitude differences in the *I* band. We also calculated the masses of each individual component and estimated orbital periods, using the available magnitude and colour relations for M dwarfs and our own  $M_J$ -spectral type and mass- $M_I$  relations. To avoid biases in our sample selection, we built a volume-limited sample of M0.0–M5.0 dwarfs that is complete up to 86% within 14 pc.

**Results.** From the 490 observed stars, we detected 80 companions in 76 systems, of which 30 are new discoveries. Another six companion candidates require additional astrometry to confirm physical binding. The multiplicity fraction in our observed sample is  $16.7 \pm 2.0\%$ . The bias-corrected multiplicity fraction in our volume-limited sample is  $19.5 \pm 2.3\%$  for angular separations of 0.2 to 5.0 arcsec (1.4–65.6 au), with a peak in the distribution of the projected physical separations at 2.5–7.5 au. For M0.0–M3.5 V primaries, our search is sensitive to mass ratios higher than 0.3 and there is a higher density of pairs with mass ratios over 0.8 compared to those at lower mass ratios. Binaries with projected physical separations shorter than 50 au also tend to be of equal mass. For 26 of our systems, we estimated orbital periods shorter than 50 a, 10 of which are presented here for the first time. We measured variations in angular separation and position angle that are due to orbital motions in 17 of these systems. The contribution of binaries and multiples with angular separations shorter than 0.2 arcsec, longer than 5.0 arcsec, and of spectroscopic binaries identified from previous searches, although not complete, may increase the multiplicity fraction of M dwarfs in our volume-limited sample to at least 36%.

**Key words.** binaries : close – stars: late-type – stars: low-mass

## 1. Introduction

The multiplicity of low-mass stars provides constraints to models of stellar and planet formation and evolution (Goodwin et al. 2007; Burgasser et al. 2007; Duchêne & Kraus 2013). M dwarfs, which have approximate masses of between 0.1 and 0.6  $M_{\odot}$ , account for two thirds of the stars in the solar neighbourhood and probably the Galaxy. However, in spite of their abundance and the increasing number of M-dwarf high-resolution imaging surveys in the past decade (Beuzit et al. 2004; Law et al. 2008; Bergfors et al. 2010; Janson et al. 2012, 2014a; Jódar et al. 2013; Bowler et al. 2015; Ward-Duong et al. 2015), the multiplicity

of M dwarfs is not yet well constrained, at least by comparison with the better determination for Sun-like stars (Duquennoy & Mayor 1991; Raghavan et al. 2010; Tokovinin 2011). Published values range between 13.6% and 42%. Thus, the binary fraction of M dwarfs seems intermediate between the one of Sun-like stars and very low mass binaries. In Table 1 we summarise the multiplicity fractions and semi-major axis coverage of some of the main multiplicity surveys carried out from F6 to T dwarfs.

The typical separation of low-mass stars in a binary system tends to decrease with the mass of the primary, which makes the detection of faint companions at resolvable separations more difficult (Jeffries & Maxted 2005; Burgasser et al. 2007; Caballero 2007; Bate 2012; Luhman 2012). In addition, the presence of a stellar companion influences planet formation (Wang et al. 2014a,b, 2015a,b). The limited number of exoplanet

<sup>★</sup> Tables A.1–A.6 are only available at the CDS via anonymous ftp to [cdsarc.u-strasbg.fr](http://cdsarc.u-strasbg.fr) (130.79.128.5) or via <http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/597/A47>

**Table 1.** Stellar multiplicity fractions.

Reference	Investigated spectral type	$d_{\text{lim}}$ [pc]	Multiplicity fraction [%]	Projected physical separation, $s$ [au]	Survey method <sup>a</sup>
Duquennoy & Mayor (1991)	F7–G9	22	~65	~0.01–225	RV, WI
Raghavan et al. (2010)	~F6–K3	25	$44 \pm 3$	~0.005–100 000	RV, AO, S, WI
Reid & Gizis (1997)	K2–M6	8	32	~0.1–1800	RV, S, WI
Leinert et al. (1997)	M0–M6	5	$26 \pm 9$	~1–100	S
Fischer & Marcy (1992)	M	20	$42 \pm 9$	0–10 000	RV, WI
Jódar et al. (2013)	K5–M4	25	$20.3^{+6.9}_{-5.2}$	~0–80	LI
Ward-Duong et al. (2015)	K7–M6	15	$23.5 \pm 3.2$	~3–10 000	AO, WI
Bergfors et al. (2010)	M0.0–M6.0	52	$32 \pm 6$	3–180	LI
Janson et al. (2012)	M0.0–M5.0	52	$27 \pm 3$	3–227	LI
Law et al. (2008)	M4.5–M6.0	$\langle 15.4 \rangle$	$13.6^{+6.5}_{-4.0}$	~0–80	LI
Siegler et al. (2005)	M6.0–M7.5	30	$9^{+4}_{-3}$	$\geq 3$	AO
Janson et al. (2014a)	M5.0–M8.0	36	21–27	~0.5–100	LI
Close et al. (2003)	M8.0–L0.5	33	$15 \pm 7$	$< 15$	AO
Bouy et al. (2003)	M7.0–L8.0	20	10–15	1–8	HST
Reid et al. (2008)	L	20	$12.5^{+5.3}_{-3.0}$	$< 3$	HST
Burgasser et al. (2003)	T	$\langle 10 \rangle$	$9^{+15}_{-4}$	1–5	HST

**Notes.** <sup>(a)</sup> AO: Adaptive optics; HST: *Hubble* Space Telescope; LI: Lucky imaging; RV: Radial velocity; S: Speckle; WI: Wide-field imaging.

hosts in binary and multiple systems (Mugrauer et al. 2007; Mugrauer & Neuhauser 2009; Ginski et al. 2015) and the relatively small number of M dwarfs with known exoplanets detected with radial-velocity and transit methods (Rivera et al. 2005; Charbonneau et al. 2009; Bonfils et al. 2013) prevents a significant statistical analysis of how stellar multiplicity at such low masses affects planet formation.

Because of their low effective temperatures, M dwarfs emit the bulk of their energy in the near-infrared. It makes them difficult to observe with the required radial-velocity precision with the current spectrographs for exoplanet hunting (e.g. HARPS at the 3.6 m ESO La Silla Telescope, HARPS-N at the 3.6 m TNG, and UVES at the 8.2 m ESO VLT), which operate in the optical. The prompt development of stable near-infrared spectrographs with wide wavelength coverage and high spectral resolution for radial-velocity surveys of M dwarfs has therefore been identified as critical by numerous decadal panels, funding agencies, and international consortia. Some noteworthy high-resolution near-infrared spectrographs currently under development are IRD at 8.2 m Subaru (Tamura et al. 2012), HPF at 9.2 m HET (Mahadevan et al. 2014), and SPIRou at 3.6 m CFHT (Donati et al. 2014). The high-resolution spectrograph CARMENES (Amado et al. 2013; Quirrenbach et al. 2014<sup>1</sup>) at 3.5 m Calar Alto covers from 520 nm to 1710 nm and has started its science survey in January 2016.

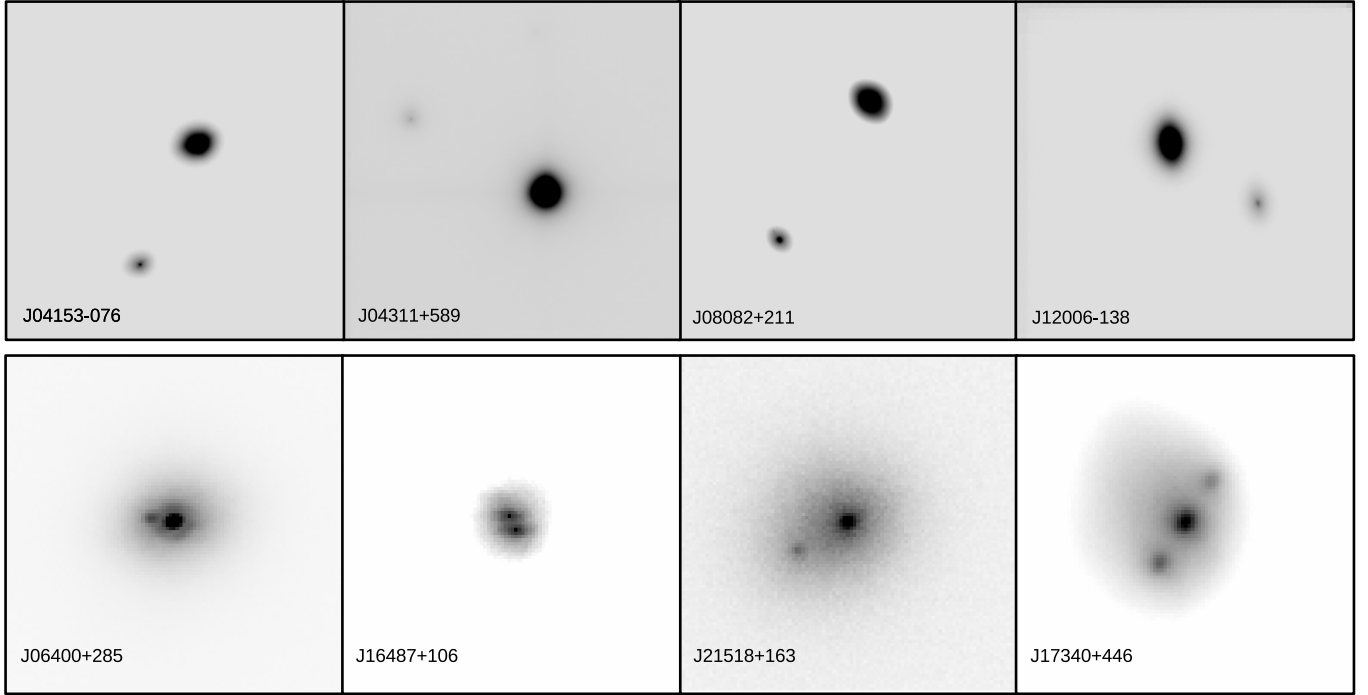
CARMENES is the name of the double-channel spectrograph (near-infrared and optical) of the Spanish-German consortium that built it, and of the science project that is being carried out during guaranteed-time observations (GTO). For at least 600 GTO clear nights in the time frame between 2016 and 2018, CARMENES will spectroscopically monitor about 300 carefully selected M dwarfs with the goal of detecting low-mass planets in their habitable zones. With a long-term  $1 \text{ m s}^{-1}$  radial-velocity precision, the consortium aims at being able to detect

$2 M_{\oplus}$  planets orbiting in the habitable zone of M5 V stars and super-Earths around earlier stars (García-Piquer et al. 2016). In addition to the detection of the individual planets themselves, the ensemble of objects will provide sufficient statistics to assess the overall distribution of planets around M dwarfs: frequency, masses, and orbital parameters.

To optimise the observational strategy of the instrument and its scientific return, the consortium has built Carmencita, the CARMENES input catalogue (Caballero et al. 2013; Quirrenbach et al. 2015; Alonso-Floriano et al. 2015). It consists of almost 2200 of the brightest M dwarfs of each spectral subtype observable from Calar Alto, from which we will select the approximately 300 single GTO stars. By single we mean stars without close visual (physically bound) or optical (unbound) stellar or substellar companions that may induce real or artificial radial-velocity variations and, therefore, contaminate the precise CARMENES measurements (Guenther & Wuchterl 2003; Ehrenreich et al. 2010; Guenther & Tal-Or 2010; Bonfils et al. 2013).

As part of our efforts to determine the multiplicity of M dwarfs and to select the best targets for radial-velocity surveys for exoplanets, we performed a high-resolution imaging search of close companions with the lucky imaging instrument FastCam at the Telescopio Carlos Sánchez, as described in this paper. Preliminary results of this work were presented as conference proceedings by Béjar et al. (2012) and Cortés-Contreras et al. (2015a,b). This paper is the second item of the series called the CARMENES input catalogue of M dwarfs. In the first paper, Alonso-Floriano et al. (2015) carried out a low-resolution optical spectroscopic analysis of a number of poorly known dwarfs to constrain their spectral types. Furthermore, this work will soon be complemented with on-going searches of unresolved spectroscopic binaries and triples identified in a large collection of high-resolution optical spectra (Montes et al. 2015; Jeffers et al., in prep.) and of wide companions to M dwarfs supported by

<sup>1</sup> <http://carmenes.caha.es>



**Fig. 1.** Selection of images of multiple systems identified by us with FastCam. North is up and east is left. The *upper row* scale is  $20 \times 20 \text{ arcsec}^2$ , that of the *lower row*  $4 \times 4 \text{ arcsec}^2$ . Images at the top were obtained with the shift & add mode, while the bottom images were obtained with the lucky image mode. The bottom right image (J17340+446) is an example of the so-called false triple effect.

virtual observatory tools (cf., Cortés-Contreras et al. 2013, 2014; Alonso-Floriano et al. 2015).

## 2. Observations

Of the almost 2200 M dwarfs currently in Carmencita, we selected 490 Carmencita targets for being observed with the FastCam lucky imager (Oscos et al. 2008) at the 1.5 m Telescopio Carlos Sánchez at the Observatorio del Teide (Tenerife, Spain). The high-resolution imager FastCam is equipped with an L3CCD Andor 512×512 detector with very low electron noise and high readout speed. It has a field of view of  $21.2 \times 21.2 \text{ arcsec}^2$  and an approximate pixel size and orientation of the detector of 0.0425 arcsec and 91.9 deg, respectively. FastCam delivers nearly diffraction-limited images, which at the Telescopio Carlos Sánchez and in the *I* band have full-width at half maxima of approximately 0.15 arcsec.

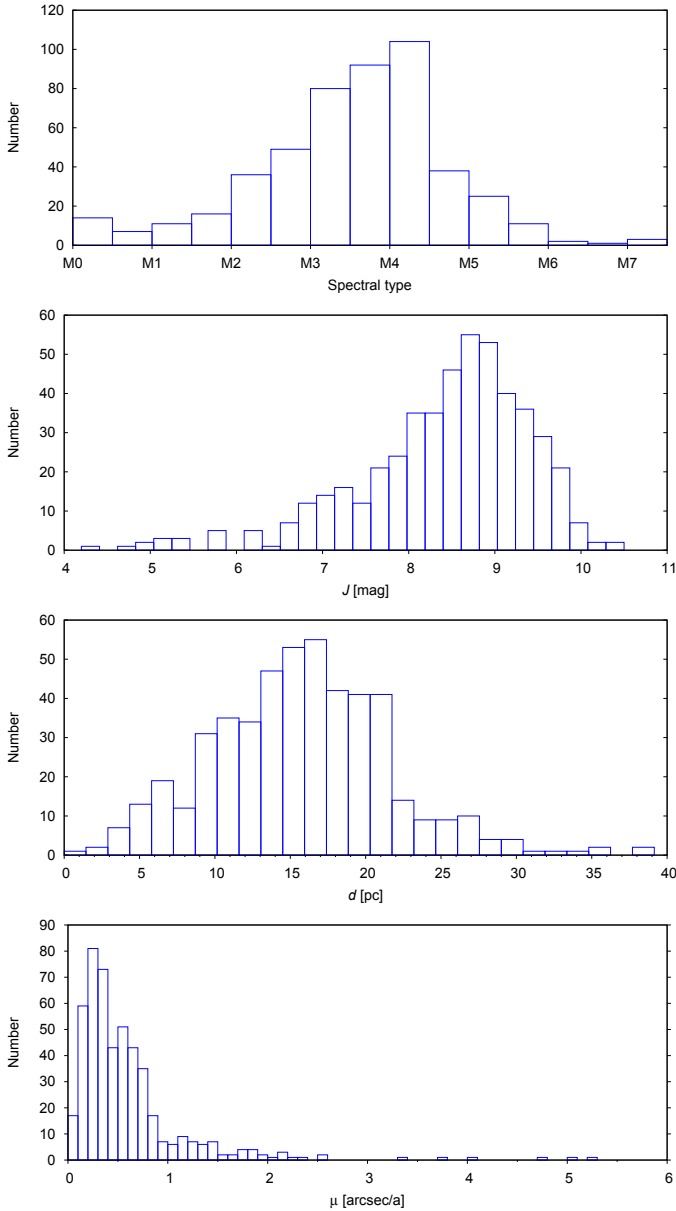
We carried out the observations during 26 nights in 15 runs from October 2011 to January 2016. For each target, we obtained typically ten blocks of 1000 frames each in the Johnson-Cousins *I* band using the electron multiplication mode. Typical frame exposure times were in the 35–50 ms range. On average, each star was imaged during 500 s in total. The typical Strehl ratio in our observations varies with the percentage of the best-quality frames chosen in the reduction process: from 0.2 for the 100% to 0.4 for the 1%. For astrometric calibration purposes, we also observed the globular cluster M3 and 18 astrometric standard binary stars from the Aitken Double Star catalogue (ADS – Aitken 1932; Scardia et al. 1995) with the same method and on several occasions.

Each frame was bias subtracted and then processed with the FastCam dedicated software developed at the Universidad Politécnica de Cartagena (see Labadie et al. 2010; Jódar et al. 2013). We ran the lucky image (on five blocks) and shift & add processing modes separately. The first allows selecting the

fraction of the best-quality frames (we chose 1%, 10%, and 50%), aligns the selected frames using the brightest speckle, and combines them, producing six final lucky images per target. The second mode aligns all the block frames and then combines them, resulting in one unique image per target. Shift & add produces deeper images than the lucky image mode, but with slightly poorer resolution. The M3 standard field was reduced only with the shift & add mode. In Fig. 1 we show a selection of the processed images at two different spatial scales.

In Table A.1, we provide the list of 490 observed M-dwarf targets with the following column information: identification number, our Carmencita identifier (Quirrenbach et al. 2015; Alonso-Floriano et al. 2015), J2000 coordinates and *J*-band magnitude from the Two-Micron All-Sky Survey (Skrutskie et al. 2006), spectral type and its reference, distance and its reference, and the FastCam observation date and exposure time. Figure 2 shows the histograms of spectral types, *J*-band magnitudes, heliocentric distances, and total proper motions of the observed sample. Spectral types range from M0.0 V to M7.0 V, *J* from 4.2 mag to 10.4 mag, distances from 1.8 pc to 39.1 pc, and proper motions from  $0.03 \text{ arcsec a}^{-1}$  to  $10.6 \text{ arcsec a}^{-1}$ . Because of their closeness, 97% of our targets have total proper motions larger than  $100 \text{ mas a}^{-1}$ .

Our sample of 490 observed Carmencita targets consisted mainly of the brightest stars in the *J* band for each spectral subtype (see Sect. 2 in Alonso-Floriano et al. 2015) that were (i) not known spectroscopic binaries; (ii) not resolved systems with visual or optical companions at angular separations smaller than 5 arcsec; and (iii) not studied with high-resolution imaging devices before the start of our observations by speckle, adaptive optics, or lucky imaging (Beuzit et al. 2004; Law et al. 2008; Bergfors et al. 2010; Jódar et al. 2013; Janson et al. 2012). Some high-resolution imaging (Janson et al. 2014a; Ansdell et al. 2015; Bowler et al. 2015; Ward-Duong et al. 2015) and spectroscopic (Bonfils et al. 2013; Llamas 2014; Schöfer 2015)



**Fig. 2.** Distributions of spectral type,  $J$ -band magnitude, distance, and proper motion of the 490 observed M-dwarf targets. The sizes of the bins follow the definitions given by Freedman & Diaconis (1981). The lowest panel does not display Barnard’s star, with  $\mu = 10.4 \text{ arcsec a}^{-1}$ .

surveys have been performed afterwards and have tabulated several objects in common with our target list. In addition, we also observed (a) some dubious or poorly investigated close multiple systems (including spectroscopic binary candidates); (b) a few stars with possible visual companions at angular separations smaller than 5 arcsec that needed confirmation or better characterisation; and (c) four known binaries with estimated orbital periods shorter than five years that were previously proposed for follow-up by Cortés-Contreras et al. (2013): J05085–181 (GJ 190), J13317+292 (DG CVn), J23174+196 (G 067–053), and J23455–161 (LP 823–004).

To confirm the physical binding of pairs (i.e. that the components share a common proper motion), we observed 54 targets more than once, and up to eight times. Accounting for the 490 M dwarfs, 18 ADS pairs and M3 calibration field, and

**Table 2.** FastCam adopted plate scale and orientation for each run night.

Observation date <sup>a</sup>	Pixel scale [mas/pix]		Orientation [deg]	
	$x$	$y$	$x$	$y$
23 Oct. 2011*	42.25	42.56	92.08	91.60
24 Oct. 2011	42.25	42.56	92.08	91.60
25 Oct. 2011	42.25	42.56	92.08	91.60
30 Jan. 2012	42.25	42.56	92.08	91.60
31 Jan. 2012	42.25	42.56	92.08	91.60
25 Mar. 2012*	42.31	42.61	91.79	91.64
26 Mar. 2012*	42.30	42.62	91.82	91.65
27 Mar. 2012	42.30	42.62	91.82	91.65
10 Jul. 2012*	42.48	42.61	92.11	91.91
11 Jul. 2012*	42.49	42.64	92.03	91.77
12 Jul. 2012*	42.32	42.54	91.96	91.99
16 Sep. 2012	42.32	42.54	91.96	91.99
17 Sep. 2012	42.32	42.54	91.96	91.99
13 Jan. 2013*	42.26	42.69	91.94	91.63
14 Jan. 2013*	42.21	42.59	91.85	91.63
28 Feb. 2014*	42.26	42.69	91.99	91.70
01 Mar. 2014	42.26	42.69	91.99	91.70
02 Mar. 2014	42.26	42.69	91.99	91.70
22 May 2014	42.26	42.69	91.99	91.70
09 Dec. 2014*	42.26	42.99	91.97	91.96
14 Apr. 2015*	42.28	42.37	92.18	91.65
15 Apr. 2015	42.28	42.37	92.18	91.65
09 Jun. 2015	42.28	42.37	92.18	91.65
29 Jul. 2015	42.28	42.37	92.18	91.65
17 Nov. 2015	42.28	42.37	92.18	91.65
07 Jan. 2016	42.28	42.37	92.18	91.65

**Notes.** <sup>(a)</sup> M3 calibration field was observed on nights marked with an asterisk.

the different epochs, we acquired 7670 images in total with FastCam.

### 3. Analysis

#### 3.1. Astrometry

The first step of the analysis was computing the pixel size and detector orientation with common IRAF tasks (Tody 1986). To do this, we determined the centroids of the brightest stars in the M3 standard field with `imcentroid`. Using the celestial coordinates in the ACS Survey of Galactic Globular Clusters (Sarajedini et al. 2007) and the pixel coordinates in our images, we then determined the transformation equations with `ccmap` by fitting to a general transformation of order two. Table 2 lists the pixel scales and orientations of the detector for each night. For nights without M3 images, we used the calibration of the closest night with computed plate solution. Pixel scale and rotation angle in the centre of the detector in the  $x$  and  $y$  axes are similar within the different campaigns with almost negligible variations from night to night. Their mean values are  $42.31 \pm 0.09 \text{ mas/pixel}$  and  $42.63 \pm 0.15 \text{ mas/pixel}$  in pixel scale and  $91.98 \pm 0.12 \text{ deg}$  and  $91.74 \pm 0.15 \text{ deg}$  in orientations of the detector in the  $x$  and  $y$  axes, respectively. The uncertainties are the standard deviations of the measurements.



To double-check that our astrometric solutions were correct, we calculated angular separations ( $\rho$ ) and position angles ( $\theta$ ) for each ADS binary. To do that, we measured the  $x$  and  $y$  positions of each star with `imcentroid`, and transformed them into equatorial coordinates using the astrometric solution of the corresponding night with `cctran`. Table A.2 shows the previously published values of  $\rho$ ,  $\theta$ , the epochs of observation and references, and our measured values in different epochs. Our errors in  $\rho$  and  $\theta$  were derived from the standard deviation of the measurements in all images within the same night and the determined astrometric solutions on different nights. In general, the measured values of  $\rho$  and  $\theta$  of the same pair on different nights were consistent within  $3\sigma$  between them and with tabulated values from recent works. In some cases, the quality of our measurements surpassed previous publications.

We carried out a visual inspection for companions to our 490 Carmencita targets and found 137 additional sources in 116 systems, for which we measured the relative positions and position angles following the same procedure as described above for the ADS binaries. In some epochs of nine stars with companions very close to the resolution limit of our images, we were unable to measure the photocentroid of both components with `imcentroid` and, hence, we used the brightest pixel to measure their positions. In these cases, the uncertainties in the determination of  $\rho$  and  $\theta$  were larger and we adopted a typical error bar of one pixel.

We classified the 137 sources into three groups: (i) 51 optical companions (i.e. unbound, Table A.3); (ii) 80 physical companions (i.e. bound, Table A.4); and (iii) six unconfirmed companions (bottom of Table A.4). For the classification, we used old photographic plate digitisations and all-sky surveys provided by the Aladin sky atlas (Bonnarel et al. 2000), previous astrometry tabulated by the Washington Double Star Catalogue (WDS, Mason et al. 2001) and/or our own multi-epoch astrometric measurements together with the target proper motions (mostly from van Leeuwen 2007 and Röser et al. 2010).

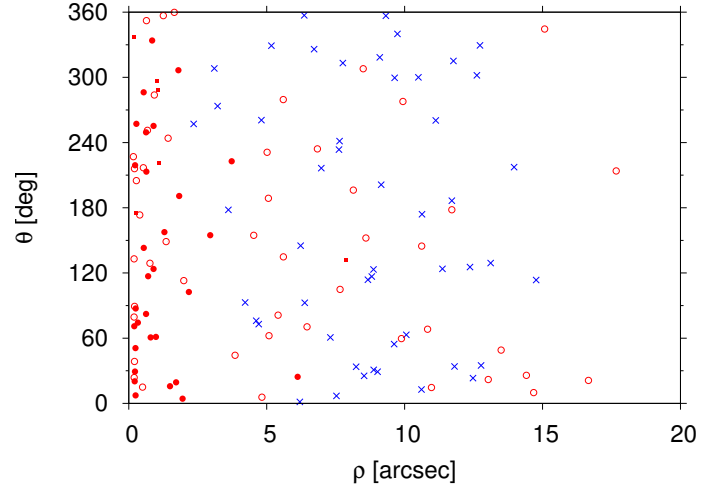
Most of the 51 optical companions are null-proper-motion point-like sources in photographic plates of the first National Geographic Society – Palomar Observatory Sky Survey in the mid-1950s. For the rest of the companions, we performed a multi-epoch analysis of their relative positions. We considered as optical (unbound) companions those that show  $\rho$  and  $\theta$  values in different epochs consistent within  $3\sigma$  with null proper motion and inconsistent by more than  $3\sigma$  with the proper motion of the M dwarfs. Otherwise, we considered them as physically bound. For five of the six unconfirmed binaries, we only had one epoch, and for the other (J07349+147), the  $\rho$  and  $\theta$  values at different epochs did not allow us to distinguish between null or common proper motion.

Figure 3 displays the measured  $\rho$  and  $\theta$  values of all the detected pairs. It shows a homogeneous distribution of the position angle of the companions, which discards possible false detections associated with, for example, optical ghosts.

### 3.2. Photometry

In Table A.4 we list magnitude differences in the  $I$  band for the 80 physical and six likely physical pairs. To measure the magnitude difference of the binaries, we performed aperture and point spread function (PSF) photometry using the `phot`, `psf`, and `allstar` routines in the `daophot` package of IRAF.

For wide enough pairs, we used the primary star PSF as a reference for the secondary. In these cases, magnitude differences from aperture photometry and PSF fitting did not differ



**Fig. 3.** Diagram of  $\theta$  vs.  $\rho$  for all the 137 measured pairs. Filled red circles are new physically bound pairs, small filled red squares are unconfirmed related pairs, open red circles are known physically bound pairs, and blue crosses are optically unrelated pairs.

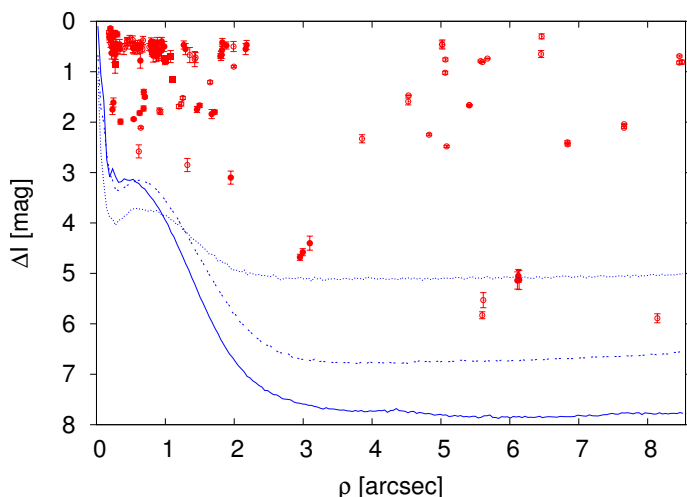
significantly. Since the PSF varies depending on the focus and sky position, for close pairs we chose the most appropriate single star observed during the same night as a reference to compute the PSF. For five pairs, we were unable to measure the  $\Delta I$  between components using PSF photometry, and we estimated it from the peak flux ratio of the PSF subtracted image, and for J23455–161, we perceived the companion and could not measure the magnitude difference.

A few close pairs showed a so-called false triple effect associated with the reduction process by the FastCam software, based on the selection of the brightest pixel. When both components are of similar brightness, this software may not distinguish between the primary and secondary and, in the process of aligning, selects the brightest pixel in one or another star, resulting in an apparent triple system. For equal brightness binaries, this may lead to a degeneracy in the determination of the position angle of 180 deg. The option `2stars` in the FastCam reduction software, which takes this ambiguity into account, solved this effect in most cases. For the rest, we determined the real flux ratio of the pair by following the procedure described by Law (2006):

$$F_R = \frac{2I_{13}}{I_{12}I_{13} + \sqrt{I_{12}^2 I_{13}^2 - 4I_{12}I_{13}}}, \quad (1)$$

where  $I_{12} = F_1/F_2$  and  $I_{13} = F_1/F_3$ , and  $F_1$ ,  $F_2$  and  $F_3$  are the fluxes of the images in the positions of the true primary, true secondary, and spurious tertiary, respectively.

In Fig. 4 we plot the measured magnitude differences in the  $I$  band and angular separations of the companions. Most of them are of similar brightness ( $\Delta I \approx 0.0$ – $1.0$  mag) and are located at angular separations smaller than 2.5 arcsec. Figure 4 also shows the contrast curves of our survey as a function of angular separation. The maximum magnitude difference in each stacked image depends on the brightness of the primary star. For this reason, we considered three different groups in our sample according to their  $I$  magnitude, from which we selected four single stars covering different spectral types to obtain a representative mean contrast curve. For each star, we estimated the detection limit as a function of the angular separation as three times the standard deviation of the number of counts in ten-pixel-wide annuli centred on the target. This detection limit was converted into



**Fig. 4.** Diagram of  $\Delta I$  vs.  $\rho$  up to 8.5 arcsec for the physical pairs. Colour and symbol code is as in Fig. 3. Solid, dashed, and dotted lines indicate the  $3\sigma$  detection limits for primaries in the magnitude ranges  $I < 10$  mag,  $10 \leq I < 11$  mag, and  $I \geq 11$  mag, respectively.

$\Delta I$  using the peak flux value of the star. The maximum magnitude difference in the detection of possible companions at angular separations between 0.2 and 1.0 arcsec varies from 3 to 4 mag and from 5 to 7 mag at separations larger than 2 arcsec, depending on the brightness of the primary star. The limiting magnitude of our survey is about  $I \approx 17$  mag, and we were able to detect all sources brighter than this limit at angular separations greater than 3 arcsec. This implies that at separations larger than 3 arcsec, the detection of companions earlier than M8 dwarfs is complete up to 40 pc, which corresponds to the entire sample, and the detection of companions earlier than M9 dwarfs is complete up to 25 pc, which is in most of our sample.

## 4. Results and discussion

### 4.1. Detected binaries

Of the 490 observed stars, we confirmed with our data 80 companions in 76 systems, of which 30 are presented here for the first time. In addition, there are also six unconfirmed binaries that need additional epochs to confirm the physical binding. The majority of the optical components of the survey were easily identified using previous available data, and most of the remaining ones were confirmed as physically bound companions using our own measurements at different epochs. Therefore, we considered the six unconfirmed binaries as very probably physically bound rather than unbound pairs. We took into account the six binaries for the determination of the multiplicity fraction.

The 86 pairs are listed in Table A.4. In the last column of the table, we include the multiplicity flag from version 1.2 of the Guide Star Catalog (Morrison et al. 2001), which is “False” for 18 of the 30 new confirmed binaries, “True” for 10 of them and has no entry for the close companions of J08082+211 and J15191–127. For the 30 new binaries, and to our knowledge, there are no other references to binarity.

Of the 80 physical companions, 48 are tabulated by WDS (second column in Table A.4), of which two were previously suggested by Behall & Harrington (1976; J05333+448) and Bowler et al. (2015; J15496+348) and confirmed here. Another two were recently presented by Ward-Duong et al. (2015; J05034+531) and Ansdell et al. (2015; J06212+442), and one

of the new companions resolved here is most likely associated with a spectroscopic binary identified by Bonfils et al. (2013; J15191–127). The remaining 29 are pairs with no previous binarity references to our knowledge.

Some of the measured companions were not detected in all epochs because of the relative motion of the components and the crossing of the companion behind or in front of the primary star (J05078+179 and J05333+448), presence of the companion near the diffraction limit (J13317+292, J15496+348, J16487+106, and J21012+332), and a focus problem (J06400+285).

### 4.2. Multiplicity fraction

Of the 490 observed M dwarfs, 408 are single and 82 (76+6) are in binary or multiple systems within the FastCam field of view. This gives a close multiplicity fraction of  $16.7 \pm 2.0\%$ , by assuming a Poissonian distribution of the errors. Nevertheless, it must not be taken as a real M-dwarf multiplicity fraction because of the selection bias of the observed sample: we did not include many stars that were previously observed in similar studies or that had known visual companions at less than 5 arcsec.

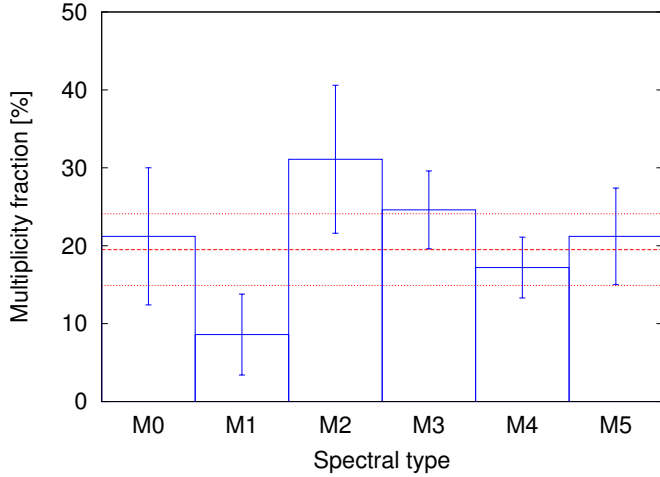
For statistical purposes, we grouped all our Carmencita (Sect. 1) and FastCam targets in a combined sample. Of the 2176 Carmencita stars, 1141 M dwarfs have been surveyed with FastCam or with high-resolution imagers with similar capabilities (Beuzit et al. 2004; Lafrenière et al. 2007; Law et al. 2008; Bergfors et al. 2010; Janson et al. 2012, 2014a; Jódar et al. 2013; Bowler et al. 2015; Ward-Duong et al. 2015). For completeness, we considered a range of angular separations from 0.2 to 5 arcsec to our targets. The lower limit was given by the FastCam spatial resolution and the upper limit by the maximum separation at which we could detect companions to at least 90% of the observed stars. Of the 1141 surveyed M dwarfs, 219 have physical companions in this interval of angular separations (55 from this work and 164 from other publications), which gives a close multiplicity fraction of  $19.2 \pm 1.4\%$ .

To avoid any selection bias and give a more reliable multiplicity fraction, we proceeded by building a volume-limited sample with a maximum distance of 14 pc and a completeness of 86%. This completeness was estimated by assuming that all M0–M5 dwarfs are known within 7 pc and that their density in the solar vicinity is constant. This third sample is composed of 425 dwarfs with spectral types between M0.0 V and M5.0 V, of which 83 have companions (either from FastCam and other works) in the range from 0.2 to 5.0 arcsec. This translates into a close multiplicity fraction of  $19.5 \pm 2.3\%$ , which is consistent within error bars with the 13.6–27% fractions obtained for M dwarfs in most surveys (Table 1), although some authors provided higher multiplicity fractions (Fischer & Marcy 1992; Bergfors et al. 2010). In Sects. 4.8 and 4.9 we estimate the contribution to the multiplicity fraction of pairs separated by less than 0.2 arcsec and more than 5 arcsec.

### 4.3. Dependence of multiplicity on spectral type

To estimate the spectral types of the individual components of the binaries, we used the  $I - J$  colours and  $M_I$  absolute magnitudes as a function of spectral type for M dwarfs from Table 3 in Kirkpatrick et al. (1994), together with the 2MASS photometry and spectral type of the pair.

For pairs resolved by 2MASS, we used these relations and values to obtain the  $I$  magnitude of the primary, and obtained the  $I$  magnitude of the secondary from our measured  $\Delta I$ . We



**Fig. 5.** Multiplicity fraction as a function of spectral type from M0.0 V to M5.0 V in the volume-limited sample with separations from 0.2 to 5.0 arcsec. Error bars are Poissonian. Horizontal dashed and dotted lines are the global multiplicity fraction and the  $\pm 2\sigma$  values.

derived the absolute  $M_I$  magnitudes through the distance modulus and inferred the spectral types of the secondaries with the  $M_I$ -spectral type relation of Kirkpatrick et al. (1994).

For pairs not resolved by 2MASS, the  $J$ -band magnitude involves the contribution of all the components in the system. In these cases, we obtained the  $I$  magnitude of the system from the  $I - J$  colours and the global spectral types of the pairs from the literature. Using the  $I$  magnitude and our measured  $\Delta I$ , we computed the individual  $I$  magnitudes. We calculated the individual  $M_I$  absolute magnitudes by applying the distance modulus, and estimated individual spectral types from the  $M_I$ -spectral type relation.

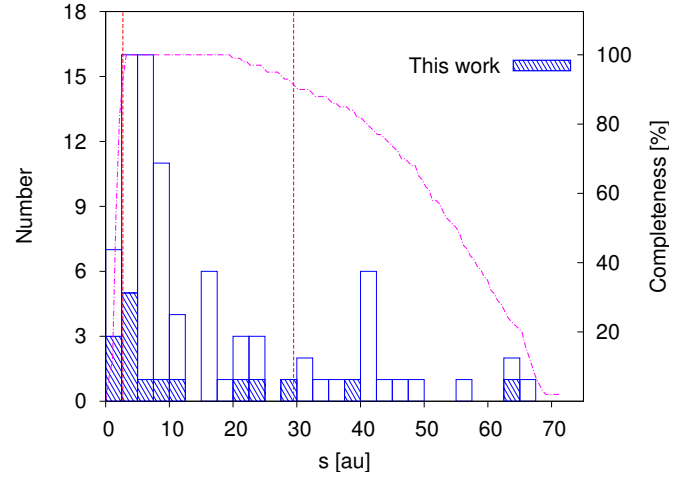
The distances in our sample come mostly from literature parallax determinations (see references in Table A.1). For stars without parallactic distance, we calculated spectro-photometric distances from our own  $M_J$ -spectral type relation. This relation was obtained from a polynomial fit using single stars with well-determined spectral types between M0 V and M6 V, parallactic distances, and 2MASS  $J$ -band photometry from the Carmencita sample, and has the form:

$$M_J = a \text{ SpT}^2 + b \text{ SpT} + c, \quad (2)$$

where  $a = 0.078 \pm 0.007 \text{ mag}$ ,  $b = 0.265 \pm 0.038 \text{ mag}$  and  $c = 5.895 \pm 0.044 \text{ mag}$ , and SpT indicates the numerical spectral subtype within the M range.

For very close binaries, spectro-photometric distances are not reliable since the 2MASS photometry and the spectral type determination do not provide the contribution of the two components separately. In these cases, in an iterative way, we estimated new individual spectro-photometric distances for the two components in the system from spectral type estimations based on the global spectral type, the  $M_I$ -spectral type relation, the individual  $I$  magnitudes, and the distance modulus. These updated distances are given in Table A.1. Given the low number of close binaries not resolved in our survey ( $\sim 10\%$ , see Sect. 4.8), we do not expect many additional unresolved components.

The individual spectral types are listed in Table A.5. SpT column indicates the combined spectral type of the system from which individual spectral types were derived. In the SpT<sub>1</sub> and SpT<sub>2</sub> columns, the spectral types indicated with capital “M” come from the literature, and with lower case “m” refer to our estimated spectral types.



**Fig. 6.** Projected physical separation distribution of the binaries in the volume-limited sample. Dashed bars represent our binaries. Vertical dashed lines mark the 90% completeness limits, and the dash-dotted curve represents the completeness as a function of the projected physical separation.

In Fig. 5 we show the dependence of the multiplicity fraction of M dwarfs on the spectral type in our volume-limited sample. The multiplicity fractions for different spectral subtypes are consistent within the error bars among them, except for M1 stars, for which it is lower. We compared this distribution with the global multiplicity fraction obtained in the previous section and performed a  $\chi^2$  test. Without the M1 contribution, the distribution is consistent with a flat distribution with a confident level of 96%.

In addition, our determined multiplicity fraction has intermediate values between Sun-like (44–65%) and very low mass stars and brown dwarfs (9–15%). This agrees with the generally accepted decreasing trend of the multiplicity fraction with decreasing mass of the primaries (Table 1).

#### 4.4. Projected physical separation distribution

To study the distribution of the binaries in the volume-limited sample, we converted angular separations ( $\rho$ ) into projected physical separations ( $s$ ) by using the small-angle approximation  $\tan \rho \approx \rho$ . Hence,  $s \approx \rho d$ . The distances  $d$  come from parallax or photometry as in Sect. 4.3.

In Fig. 6 we show the projected physical separation distribution of the binaries in the volume-limited sample. We also represent the completeness of the volume-limited sample as a function of projected physical separation, and draw the completeness limits with a confidence level of 90%, which correspond to the  $s$  interval between 2.6 and 29.5 au. We estimated these values as those separations at which we are able to detect companions in 90% of the sample.

The projected physical separations of the observed pairs in the volume-limited sample range from 1.4 to 65.6 au and their distribution peaks at 2.5–7.5 au. This is consistent with the values of 5–10 au found by Jódar et al. (2013) for M0–M4 dwarfs, and of  $\sim 6$  au found by Janson et al. (2014) for M3–M8 dwarfs and Ward-Duong et al. (2015) for M0–M6 dwarfs. However, these values are lower than those found for Sun-like stars (Duquennoy & Mayor 1991; Raghavan et al. 2010) and more similar to those found for ultracool dwarfs (4–6 au for M8.0–L0.5, Close et al. 2003; 2–4 au for M7.0–L8.0, Bouy et al. 2003;  $< 3$  au for L dwarfs, Reid et al. 2008).

**Table 3.** Target members of stellar kinematic groups.

Karmn	Name	Moving group	Ref. <sup>a</sup>	Assumed age [Ma]	Ref. <sup>b</sup>
J01221+221	G 034–023	Young disc	Abe14	≥300	This work
J04153–076	<i>o</i> <sup>02</sup> Eri C	$\beta$ Pic	AF15	~20	Bell15
J05019+099	LP 476–207	$\beta$ Pic	AF15	~20	Bell15
J05068–215E	BD–21 1074 A	$\beta$ Pic	AF15	~20	Bell15
J05068–215W	BD–21 1074 BC	$\beta$ Pic	AF15	~20	Bell15
J05103+488	G 096–021 AB	IC 2391?	This work	~50	Barr04
J10028+484	G 195–055	Local Association?	This work	~100	Bas96
J10196+198	BD+20 2465	Castor	Cab10	≥300	Barr98, Mam13
J12123+544S	BD+55 1519 A	UMa	Mon01	≥300	Gia79, SM93
J12123+544N	BD+55 1519 B	UMa	Mon01	≥300	Gia79, SM93
J13317+292	DG CVn AB	Columba/Carina	Ried14	~40	Bell15
J18548+109	V 1436 Aql B	Castor	Cab10	≥300	Barr98, Mam13
J23293+414S	G 190–027	Local Association	Klu14	~100	Bas96
J23293+414N	G 190–028	Local Association	Klu14	~100	Bas96
J23318+199 E	EQ Peg Aab	Castor	Cab10	≥300	Barr98, Mam13
J23318+199 W	EQ Peg Bab	Castor	Cab10	≥300	Barr98, Mam13

**Notes.** <sup>(a)</sup> Abe14: Aberasturi et al. (2014); AF15: Alonso-Floriano et al. (2015); Cab10: Caballero (2010); Klu14: Klutsch et al. (2014); Mon01: Montes et al. (2001); Ried14: Riedel et al. (2014). <sup>(b)</sup> Barr04: Barrado y Navascués et al. (2004); Barr98: Barrado y Navascués (1998); Bas96: Basri et al. (1996); Bell15: Bell et al. (2015); Gia79: Giannuzzi (1979); Mam13: Mamajek et al. (2013); SM93: Soderblom & Mayor (1993).

Within the physical separation completeness range from 2.6 to 29.5 au, there are 61 M dwarfs with low-mass companions in our volume-limited sample. This translates into a multiplicity fraction of  $14.4 \pm 2.0\%$ , which is lower than the fraction derived in Sect. 4.2 as a result of the missing systems at larger separations (see Fig. 6).

#### 4.5. Masses

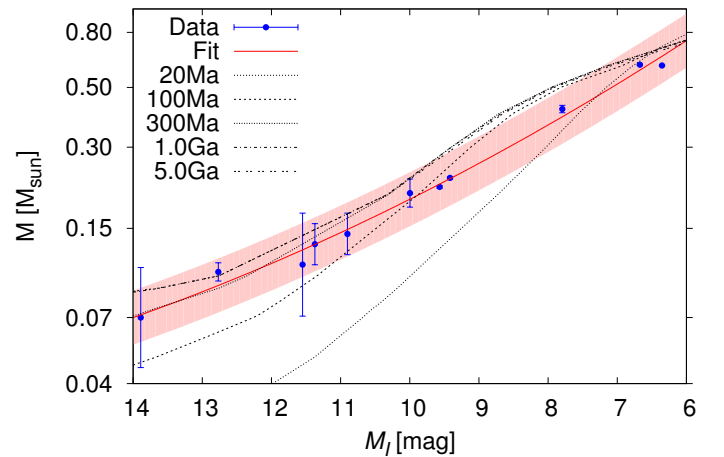
We derived masses from our own mass-luminosity relation in the Johnson-Cousins *I* band. To our knowledge, there is no published mass-luminosity relation employing this band. We collected dynamical masses and *I*-band magnitudes of eleven low-mass stars from different works (Delfosse et al. 2000; Henry 2004; Reid et al. 2004; Tokovinin 2008) and obtained an  $M_I$ - $M$  relation using a parabolic fit of the form:

$$\log M = a M_I^2 + b M_I + c, \quad (3)$$

where  $M$  is the mass,  $M_I$  is the absolute *I*-band magnitude,  $a = 0.005 \pm 0.002 \text{ mag}^{-2}$ ,  $b = -0.222 \pm 0.037 \text{ mag}^{-1}$ , and  $c = 1.035 \pm 0.180$ . This relation is valid for main-sequence stars in the  $M_I$  interval between 6 and 14 mag, which corresponds to ~M0–M8 spectral types. Figure 7 shows the data taken from the literature, the corresponding best fit, and the comparison with BT-Settl evolutionary models from the Lyon group (Baraffe et al. 2015).

In some of our detected pairs, one or both components are also spectroscopic binaries (see Table 6). For these we estimated individual masses assuming equally bright components.

For main-sequence stars, the luminosity and effective temperatures are unambiguously related to the mass, and thus, the relation in Eq. (3) is only valid for stars older than ~300 Ma, as



**Fig. 7.** Mass  $M$  vs. Absolute magnitude  $M_I$ . Blue points represent the dynamical masses and absolute magnitudes taken from the literature. The red solid line and shadowed area represent the best-fit  $\pm 3\sigma$ . Different dashed lines display the BT-Settl evolutionary models at 20 Ma, 100 Ma, 300 Ma, 1 Ga, and 5 Ga.

inferred from Fig. 7. For stars younger than 300 Ma, the mass-luminosity relation strongly depends on the age. We searched for young stars in our sample by collecting radial velocities from the literature (Caballero et al., in prep.) and computing UVW Galactocentric space velocities as in Montes et al. (2001) for 452 of the 490 observed stars (there are 38 stars without radial velocities). Of these, 155 have *U* and *V* velocity components inside or near the boundaries that delineate the young-disc population (Montes et al. 2016). In total, 42 stars of our 82 multiple systems are candidate members in young stellar kinematic groups. We checked



the literature and found that 26 of the 42 are relatively old interloper stars that do not show any youth feature or have been poorly investigated. The remaining 16 stars are confirmed members of stellar kinematic groups or the young-disc population. Their associations and ages are listed in Table 3.

Since  $I - J$  colours of young stars and field stars do not show significant differences (see Bihain et al. 2010; Peña-Ramírez et al. 2016), we applied the colour-spectral-type relation from Kirkpatrick et al. (1994) to derive the individual  $I$  magnitudes of these stars as explained in Sect. 4.3. We considered Castor, Ursa Majoris, and young-disc members old enough to be main-sequence stars, and thus, to apply our mass- $M_I$  relation with confidence. For these calculations, we assumed the ages given in Table 3. These stars appear in italics in Table A.5. The candidate pair to IC 2391 does not have a parallactic distance. Hence, we estimated its mass from the  $I - J$  colours and the BT-Settl evolutionary models from the Lyon group (Baraffe et al. 2015). We also applied these models to derive masses from the individual  $I$  magnitudes and parallactic distances for  $\beta$  Pic, Columba/Carina and Local Association members.

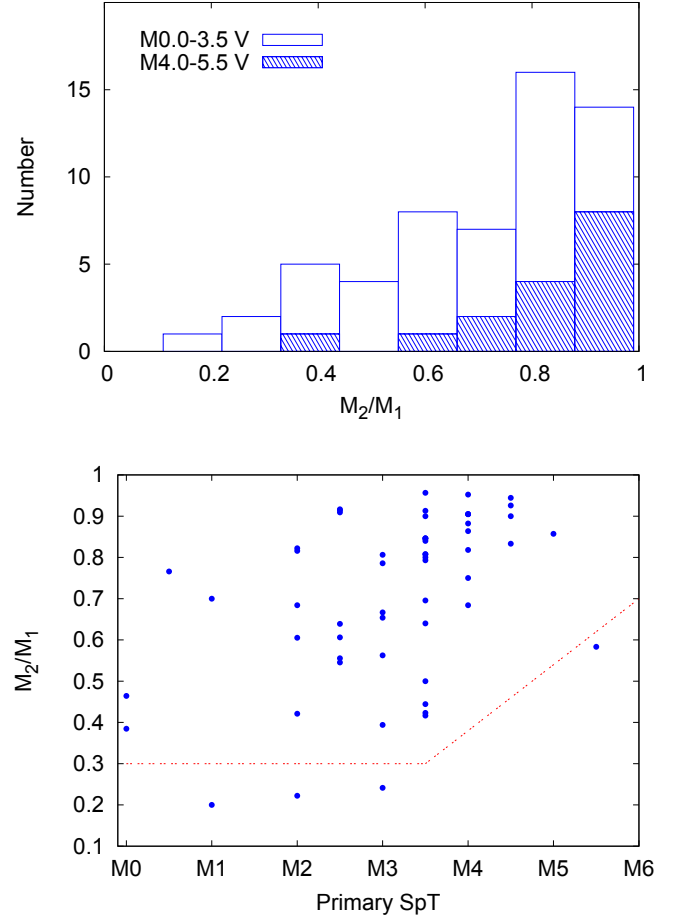
Table A.5 lists the inferred  $I$  magnitudes (Sect. 4.3) and mass values of the components of 76 of our systems. All of the detected companions have absolute magnitudes brighter than 14 mag, the lowest limit of our empirical mass-magnitude relation, which corresponds to masses close to the hydrogen-burning limit ( $\sim 0.07 M_\odot$ ). The only exception is the unconfirmed companion of J04352–161, which has an absolute magnitude fainter than 14 mag, and we were unable to determine its mass with the method explained before.

#### 4.6. Mass ratios

The upper panel in Fig. 8 shows the mass ratio ( $M_2/M_1$ ) histogram of our binaries in Table A.5. This global distribution slightly increases towards higher mass ratios and has its maximum above 0.8. The slightly lower number of equal-mass pairs with mass ratios near unity is not significant and could be related to the effect of the reduction process using the brightest pixel, which artificially sharpens the PSF of the primary with respect to the PSF of the secondary, and may produce a lower flux ratio than expected. This distribution is also affected by our sensitivity limit. While in spectral types earlier than M3.5 (i.e. more massive stars) our search of companions is complete for mass ratios greater than 0.3, in later spectral types (i.e. less massive stars) the search is complete for mass ratios greater than 0.35–0.60.

Empty and dashed bars represent the mass ratio distributions of M0.0–M3.5 and M4.0–M5.5 primaries, respectively. The distribution of the former shows the same trend as the global distribution, with a peak around 0.8–0.9. For the latter, the distribution increases towards higher ratios. As explained before, this might be due to our observational bias.

The high occurrence of binaries with mass ratios above 0.8 can also be seen in the lower panel in Fig. 8, which represents the spectral type of the primary versus the mass ratio. For later spectral types, our detected binaries also tend to have similar masses. This may be due to the lack of sensitivity to lower mass ratios at later spectral types. The distribution differs with the more homogeneous mass ratio distributions observed by Janson et al. (2012, 2014). The number of binaries with mass ratios closer to unity (i.e. similar masses) for M0.0–M3.5 contrasts with the relatively low numbers presented in Bergfors et al. (2010) in this range, but is more similar to their distribution for later M4.0–M5.5 spectral types.



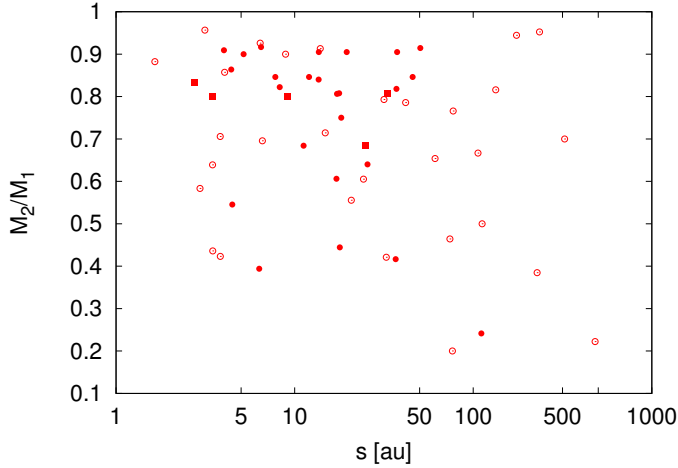
**Fig. 8.** *Top panel:* mass ratio distribution of our binaries. Empty and dashed bars separate the mass ratio distribution of M0.0–M3.5 and M4.0–M5.5 dwarfs. *Bottom panel:* mass ratio of the pairs vs. spectral type of the primary. The red dashed line represents the mass ratio completeness limits. The standard error of the mean mass ratio is 0.03 and the error bar is  $\pm 0.5$  in spectral type.

Figure 9 displays the occurrence of mass ratios with physical separations. Pairs with separations shorter than 50 au tend to have mass ratios over 0.8, while pairs at larger separations present a more homogeneous distribution.

Similar studies also show this observed trend in the relation between separation of the components and mass ratio: near equal-mass pairs (mass ratios  $\geq 0.8$ ) are found at smaller separations. Moreover, the lower the mass of the primary, the higher the mass ratio and the closer the semi-major axis at which companions are found (Jódar et al. 2013; Janson et al. 2012, 2014). The closer distance to the Sun of our sample compared to the samples of Bergfors et al. (2010) and Janson et al. (2014), who investigated the mass ratio at larger separations, may explain the difference with our results in the mass ratio distribution. However, Monte Carlo simulations of Sun-like stars and M-dwarf surveys from Duquennoy & Mayor (1991) and Raghavan et al. (2010), and Fischer & Marcy (1992) and Janson et al. (2012), respectively, suggest that the mass ratio distributions could be independent of the separation and dynamical evolution (Reggiani & Meyer 2011, 2013).

#### 4.7. Periods and orbital motion

We derived periods for 70 systems with Kepler’s third law, the masses of the components, and the maximum projected physical



**Fig. 9.** Mass ratio vs. projected physical separation. Colour and symbol code is as in Fig. 3.

separations (Sect. 4.4). Since these measures are a lower limit estimate to the semi-major axis, the periods given in Table A.5 should be also considered as a lower limit.

In total, 26 systems have periods shorter than 50 a, of which 13 are known bound systems, 10 are newly discovered binaries, and three are the unconfirmed pairs J01221+221, J07349+147, and J10028+484.

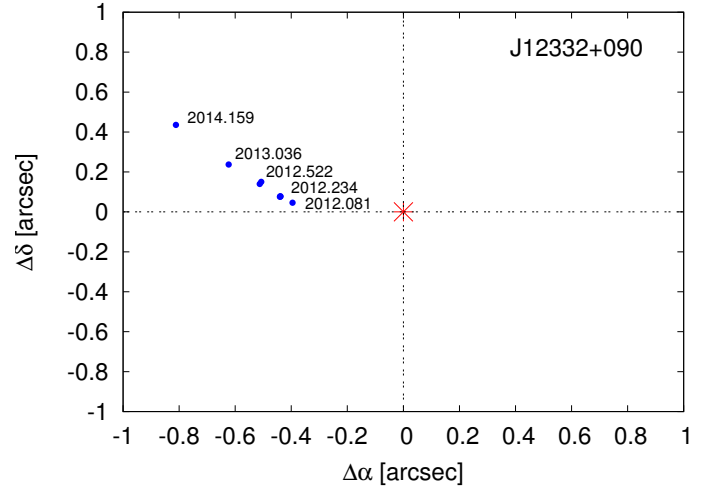
Of the 26 systems, we consider four triple systems here: J05078+179, J08082+211, and J16554–083S, which are formed by a spectroscopic binary plus a third resolved component, and J23293+414S, for which we resolved the three components of the system. In addition, the “triples” J08082+211 and J16654–083S belong to a hierarchical quadruple and quintuple system, respectively, with the fourth and fifth components outside the field of view of FastCam (Sect. 4.9).

Several systems were observed repeatedly during the programme, which allowed us to perform a multi-epoch analysis. Some of them showed appreciable variation of angular separation and position angle in different epochs of our data. When these variations were larger than  $3\sigma$  with respect to constant values of  $\rho$  and  $\theta$  and were consistent with an orbital trajectory, we considered that the orbital motion of the pair was detected. Because of the large uncertainties, the variations of  $\rho$  and  $\theta$  of the pairs J05333+448, J08066+558, and J20407+199 lie within  $3\sigma$  and therefore they do not fulfil our criterion, but they show appreciable variations that are probably related to the orbital motion. However, the time baseline is not long enough to provide a precise estimate of the orbital parameters of the systems.

Table 4 lists these 16 systems, of which 13 are new. We tabulate the WDS discoverer code of the previously known pairs, the number of used epochs, the time interval between the first and last measured epoch, and the estimated periods. We show an example of one of these binaries (J12332+090) in Fig. 10.

#### 4.8. Known close and spectroscopic binaries (not detected in our search)

In the observed sample there were also previously known pairs that we were unable to resolve because of the small separation of the components ( $\rho \lesssim 0.2$  arcsec) and/or the faintness of the companion. These pairs are listed in Table 5. In addition, there were also previously known spectroscopic binaries, taken into account for the period estimation of our detected binaries in Table A.5. They are listed in Table 6.



**Fig. 10.** Orbital variation of the pair J12332+090 from our FastCam data. The asterisk marks the position of the primary. Five of the eight epochs are labelled.

**Table 4.** Systems with measurable orbital motion.

Karmn	WDS	Epochs	$\Delta t$ [a]	$P$ [a]
J02518+294	...	3	4.2	130
J05068–215W	DON 93	3	1.2	62
J05078+179	...	2	1.1	50
J05333+448	BH 76	6	2.3	8.3
J06400+285	...	3	1.0	20
J08066+558	...	4	3.2	26
J08082+211	...	2	3.8	33
J08595+537	...	2	1.1	19
J11355+389	...	5	3.0	31
J11521+039	...	2	1.1	15
J12332+090	REU 1	8	2.1	16
J13180+022	...	4	3.1	73
J14210+275	...	2	3.0	97
J16487+106	...	3	1.1	10
J17530+169	...	5	3.0	110
J20407+199	RAO 23	2	2.8	8.4
J21518+136	...	3	3.0	66

The close multiplicity fraction of  $19.5 \pm 2.3\%$  given in Sect. 4.2 is a lower limit of the total multiplicity fraction of M dwarfs, since it only includes physical companions in the interval of angular separations between 0.2 and 5.0 arcsec. Although studies of spectroscopic binaries and very close binaries ( $\rho < 0.2$  arcsec) are not complete, we know from the literature that we are missing 47 very close additional binaries in this range in our volume-limited sample (e.g. Delfosse et al. 2013; Schöfer et al. 2015; Tokovinin et al. 2015). This number is consistent with the fractional incidence of eclipsing binaries obtained from surveys like *Kepler* (Shan et al. 2015), and increases the given binary fraction by 11%. Hence the multiplicity fraction at separations smaller than 5 arcsec would be at least  $\sim 30\%$ .

#### 4.9. Known companions at separations larger than 5 arcsec

Many of our FastCam stars have stellar or substellar companions outside the field of view of the instrument or at angular separations larger than the 5.0 arcsec cut-off defined for statistical

**Table 5.** Astrometric properties of previously known imaging companions at  $\rho < 5$  arcsec not resolved or detected in our data.

Karmn	WDS	Discoverer code	$\rho$ [arcsec]	$\theta$ [deg]	Epoch [a]	Ref. <sup>a</sup>	$\Delta\text{mag}$ (band) [mag]
J00088+208	00089+2050	BEU 1	0.133	271.9	2012.02	Jan14a	1.59 ( <i>i'</i> )
J05085-181	05086-1810	WSI 72	0.07	44.4	2011.04	WD15	0.1 ( <i>K<sub>s</sub></i> )
J04311+589	...	...	0.07	...	1965.702	Str77	0.5 ( <i>V</i> )
J06523-051	06523-0510	WSI 125	0.18	149.6	2010.068	Mas01	0.5 ( <i>o</i> )
J07307+481	...	...	0.054	...	1960.60	Harr81	...
J09177+462	09177+4612	JNN 68	0.204	37.5	2011.073	Bow15	0.102 ( <i>K<sub>s</sub></i> )
J10513+361	10513+3607	BWL 26	0.206	119.60	2012.357	Bow15	3.3 ( <i>H</i> )
J12290+417	12290+4144	BWL 31	0.0503	255.5	2011.469	Bow15	0.647 ( <i>H</i> )
J16241+483	16240+4822	HEN 1	0.1387	295.4	2006.62	Mar07	2.781 ( <i>H<sub>cont</sub></i> )
J16354+350 <sup>b</sup>	16355+3501	BWL 44	0.092	25.62	2011.469	Bow15	0.406 ( <i>H</i> )
J17177+116 <sup>c</sup>	...	...	...	...	1977	Chr78	...
J18387-144	18387-1429	HDS 2641	0.107	358	1991	DN00	0.04 ( <i>H<sub>p</sub></i> )
J19122+028	19121+0254	AST 1	0.16	319.7	2007.36	WD15	0.80 ( <i>H</i> )
J20298+096 <sup>d</sup>	20298+0941	AST 2	0.160	89.1	2012.66	Jan14a	2.72 ( <i>z'</i> )
J20433+553	20433+5521	LLO 1	0.854	20.2	2007.66	Ire08	5.06 ( <i>H</i> )
J21013+332	21013+3314	JNN 288	0.142	34.0	2012.01	Jan14a	1.07 ( <i>i'</i> )
J21160+298E	21161+2951	BWL 56	0.0543	354.6	2011.47	Bow15	0.37 ( <i>H</i> )
J21313-097	21313-0947	BLA 9	0.16	128.2	2005.33	WD15	1.12 ( <i>H</i> )
J23174+196	23175+1937	BEU 23	0.145	220.2	2012.65	Jan14b	1.17 ( <i>J</i> )

**Notes.** <sup>(a)</sup> Bow15: Bowler et al. (2015); Chr78: Christy (1978); DN00: Dommanget & Nys (2000); Harr81: Harrington et al. (1981); Ire08: Ireland et al. (2008); Jan14a: Janson et al. (2014a); Jan14b: Janson et al. (2014b); Mar07: Martinache et al. (2007); Mas01: Mason et al. (2001); Str77: Strand (1977); Tok15: Tokovinin et al. (2015); WD15: Ward-Duong et al. (2015). <sup>(b)</sup> The BWL 44 companion at 2.2 arcsec is optical. <sup>(c)</sup> Astrometric perturbation with a 10 a period estimation in Chr78. <sup>(d)</sup> Spectroscopic binary identified by Benedict et al. (2000) and resolved by Janson et al. (2014a) for the first time.

purposes. We compiled the multiplicity information of all of them using our observations and the WDS catalogue. Of the wide binaries present in the WDS catalogue, nearly 60% come from the Luyten Double Star Catalogue (Luyten 1997) and the Lowell Proper Motion Survey (Giclas et al. 1971). In Table A.6, we list for each wide system the WDS discoverer code, names, spectral types, and angular separation.

As a summary, of the 490 observed stars, 50 are M-dwarf primaries with M-type wide companions, four with white dwarf companions, one with an L-dwarf, and three with a T-dwarf secondary. In addition, 11 M secondaries have F (2), G (3), K (4) or white dwarf (3) primaries. Five tertiary M dwarfs are in triple systems involving K+M (1), G+K (1), K+DA (1) or K+K (2) primaries.

In our volume-limited sample are 25 M dwarf primaries with wide M, L, or T dwarf secondaries at separations larger than 5 arcsec. Although our search at wide separations is not complete, since we carried out a compilation from different studies in the literature, we estimated an increment in the multiplicity fraction of 6% (25 systems out of 425 M dwarfs in our volume-limited sample), which added to the percentage estimated for pairs at separations closer than 0.2 arcsec, and spectroscopic binaries would translate into a minimum multiplicity fraction at all separations of  $\sim 36\%$ .

## 5. Summary

We obtained high-resolution images in the *I* band of 490 M dwarfs of the CARMENES input catalogue (Carmencita) with the lucky imaging instrument FastCam at the 1.5 m Telescopio Carlos Sánchez.

Among the 490 observed M dwarfs, we identified 80 physically bound companions in 76 systems, of which 30 are presented here for the first time, plus six unconfirmed companions. For all of them, we measured angular separations, position angles, and *I*-band magnitude differences. From the  $\Delta I$  differences, together with 2MASS photometry, spectral type, and colour-magnitude relations for field M dwarfs, we estimated individual *I*-band magnitudes and spectral types of each component. We also derived individual masses  $M$  and estimated orbital periods for these pairs from our own  $M$ – $M_I$  relation. For these calculations, we used parallactic distances. When not available, we derived spectro-photometric distances from our determined  $M_J$ -spectral type relation.

For our observed sample, we determined a multiplicity fraction of  $16.7 \pm 2.0\%$ . However, our sample has a strong selection bias because we discarded M stars with previously known companions at separations smaller than 5 arcsec. To obtain an unbiased multiplicity fraction, we built a volume-limited sample of Carmencita stars observed with FastCam and similar high-resolution imagers. It contains 425 M0–5 dwarfs and is complete up to 86% within 14 pc. For this sample, we derived a multiplicity fraction of  $19.5 \pm 2.3\%$  in the completeness range of angular separations between 0.2 and 5.0 arcsec, which agrees with previously reported values (Leinert et al. 1997; Janson et al. 2012, 2014; Jódar et al. 2013; Ward-Duong et al. 2015). The multiplicity fraction is consistent with a flat distribution from M0 V to M5 V within Poissonian error bars, and has intermediate values between solar-type stars and very low mass stars and brown dwarfs in accordance with the decreasing tendency observed towards lower masses.

The distribution of the number of pairs as a function of projected physical separation has a maximum between 2.5 and

**Table 6.** Known spectroscopic binaries in the observed sample.

Karmn	Spectroscopic binarity	Ref. <sup>a</sup>
J03346–048	SB3	Llam14
J03526+170	SB2	Bon13
J04252+080S	SB2	Llam14
J04352–161	SB2	RB09
J04488+100	SB2	Jeff16
J05019+099	SB2	Del99
J05032+213	SB2	Jeff16
J05078+179	SB1	Jeff16
J05342+103S <sup>b</sup>	SB	Rein12
J05466+441	SB2	Jeff16
J07418+050	SB2	Llam14
J08082+211	SB2	Shk10
J09011+019	SB2	Jeff16
J09120+279	SB2	Jeff16
J09143+526	SB1	Jeff16
J11036+136	SB1	Jeff16
J12142+006	SB2	Bon13
J12191+318	SB2	Jeff16
J12290+417	SB2	Jeff16
J14171+088	SB2	Jeff16
J14368+583	SB2	Jeff16
J15191–127	SB	Bon13
J16255+260	SB2	Jeff16
J16487+106	SB2	Jeff16
J16554–083S	SB	Pett84
J18411+247S	SB2	GR96
J19354+377	SB1	Jeff16
J20433+553	SB2	Ire08
J20445+089N <sup>c</sup>	SB1	Jeff16
J23096–019	SB2	Jeff16
J23174+382	SB2	Jeff16
J23318+199E	SB1	Del99
J23318+199W	SB1	Del99
J23573–129W	SB2	Jeff16

**Notes.** <sup>(a)</sup> Bon13: Bonfils et al. (2013); Del99: Delfosse et al. (1999); GR96: Gizis & Reid (1996); Ire08: Ireland et al. (2008); Jeff16: Jeffers et al. (in prep.); Llam14: Llamas 2014; Pett84: Petersen et al. (1984); RB09: Reiners & Basri (2009); Rein12: Reiners et al. (2012); Shk10: Shkolnik et al. (2010). <sup>(b)</sup> From the spectral types and magnitude differences of the components, we infer that the spectroscopic binary is the B companion. <sup>(c)</sup> Equal-brightness close binary previously suggested by Cortés-Contreras et al. (2014).

7.5 au and decreases at wider separations. The pairs with projected physical separations smaller than 50 au tend to have mass ratios higher than 0.8, while for larger separations this distribution is more uniform.

We estimated that 26 of our systems have orbital periods shorter than 50 a, of which 10 are newly discovered systems. In 17 of them, we were able to detect orbital variations within our own multi-epoch measurements. These systems are especially interesting for future astrometric follow-up for determining their orbital solutions and measuring dynamical masses.

For our volume-limited sample, we also collected from the literature the physically bound companions at separations closer than 0.2 arcsec and larger than 5 arcsec, and unresolved spectroscopic binaries. The addition of these systems may increase the multiplicity fraction derived in this work to at least 36%,

a value consistent with the  $42 \pm 9\%$  obtained by Fischer & Marcy (1992). Nevertheless, the sample is not complete at separations beyond the completeness limit of our survey (0.2–5.0 arcsec) and, hence, this value must only be considered as a rough estimation.

Finally, we provided a complete sample of multiple M dwarfs useful for studying the effect of low-mass stellar multiplicity on planet formation with the help of CARMENES and other near-infrared high-resolution spectrographs.

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